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### Enhancement in treatment planning for magnetic nanoparticle hyperthermia: Optimization of the heat absorption pattern

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## Enhancement in treatment planning for magnetic nanoparticle hyperthermia: Optimization of the heat absorption pattern

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### Abstract

In clinical applications of magnetic nanoparticle hyperthermia for cancer treatment it is very important to ensure a maximum damage to the tumor while protecting the normal tissue. The resultant heating pattern by the nanoparticle distribution in tumor is closely related to the injection parameters. In this study we develop an optimization algorithm to inversely determine the optimum heating patterns induced by multiple nanoparticle injections in tumor models with irregular geometries. The injection site locations, thermal properties of tumor and tissue, and local blood perfusion rates are used as inputs to the algorithm to determine the optimum parameters of the heat sources for all nanoparticle injection sites. The design objective is to elevate the temperature of at least 90% of the tumor above 43°C, and to ensure only less than 10% of the normal tissue is heated to temperatures of 43°C or higher. The efficiency, flexibility and capability of this approach have been demonstrated in a case study of two tumors with simple or complicated geometry. An extensive experimental database should be developed in the future to relate the optimized heating pattern parameters found in this study to their appropriate nanoparticle concentration, injection amount, and injection rate. We believe that the optimization algorithm developed in this study can be used as a guideline for physicians to design an optimal treatment plan in magnetic nanoparticle hyperthermia.

**Keywords:** *Nanoparticle hyperthermia, optimization, temperature, injection site, SAR, tumor, cancer treatment*

### Introduction

Although magnetic nanoparticle hyperthermia was developed in the 1950s [1], it is still at an early stage of development in the treatment for tumors. In this method, superparamagnetic magnetic nanoparticles can deliver adequate heating to irregular and/or deep-seated tumors when exposed to relatively low magnetic field at low frequency. The heat generated by these particles when exposed to an external alternating magnetic field is mainly due to the Néel relaxation and/or Brownian motion of the particles [2–3]. Previous studies [4] showed that iron oxides magnetite  $\text{Fe}_3\text{O}_4$  and maghemite  $\gamma\text{-Fe}_2\text{O}_3$  nanoparticles are biocompatible in human tissue [5]. Different methods exist to deliver nanoparticles to the tumor: either systemically if the blood vessels of the diseased organ are well known,

or by directly injecting the nanoparticles into the extracellular space of the tumor. The direct injection method is the focus of this study and its advantage is that multiple-site injections can be exploited to cover the entire target region in the case of an irregularly shaped tumor. The success of this approach to kill tumor relies on its ability to elevate temperatures above a certain threshold in the entire tumor with minimal thermal damage to the surrounding healthy tissue.

This approach has been used to treat tumors in liver [6] and breast [4, 7]. The clinical study by Johannsen and Jordan [8–10] focused on testing magnetic fluid hyperthermia on prostate cancer. In the studies of Johannsen and Jordan, the histological analysis of the cancerous tissues showed a partial necrosis of the cells after the treatment. These experiments have demonstrated the feasibility

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of elevating the tumor temperature to a desired level; however, in some tumor regions, usually at the tumor periphery, under-dosage of heating (temperature increases lower than a critical value) was observed.

The achieved distribution of the temperature elevation and the heating duration are the main factors that affect the direct cell-killing and anti-cancer immune response phenomena. The temperature distribution is determined by heat conduction in tissue and distribution of the heat generation. Typically, heat generation in the hyperthermia process is described by a quantity SAR (specific absorption rate). Previous investigations have demonstrated that particle size, particle coating, and magnetic field strength and frequency determine its heating capacity defined as the specific loss power (SLP) [2, 3, 11, 12]. However, for a given type of the particles and the magnetic field strength used in a treatment procedure, the spatial distribution of the particles dispersed in tissue dominates the distribution of the SAR and the corresponding spatial temperature elevation. A recent experimental study in a tissue-equivalent agarose gel by our group has shown that the particles were confined in the vicinity of the injection site and that their concentration was not uniform after the injection. The particle deposition was greatly affected by the injection rate and amount [13]. An *in vivo* experimental study performed on rat hind limbs [14] showed that the SAR around the injection site can be approximated by Gaussian distribution  $SAR = A \cdot e^{-r^2/r_0^2}$ , where  $r$  is the spatial distance from the injection site,  $A$  represents its maximum value and  $r_0$  represents how far from the injection site the heating is affecting the tissue. A spherical SAR distribution is possible when a very low injection rate is applied [13, 14]. The SAR distribution was determined based on measured temperature elevations in the rat hind limb. It was found that the injection amount affected the parameters of the SAR distribution [14]. The knowledge of a SAR expression allows us to determine the temperature distribution in the vicinity of an injection site at a specific injection rate, amount, and initial ferrofluid concentration. A single injection usually results in a nearly spherical-shaped region with thermal damage. For large-sized and irregular-shaped tumors, it is imperative to employ a multi-site injection technique so that a sufficiently high temperature elevation within the entire tumor can be achieved. The main challenge in this type of clinical applications is how to obtain an optimized design of the locations of the injection sites and the ferrofluid injection rate, amount and initial ferrofluid concentration at each site to ensure the maximum damage to tumor and minimize the temperature elevation in the surrounding healthy tissue.

The goal of any optimization problem is to minimize or maximize a so-called ‘objective function’ as a function of a set of parameters. In hyperthermia treatments of tumors, the region in a tumor and its boundary having a temperature above a certain threshold should be maximized while the temperature elevations in the surrounding normal tissue have to be minimized. The objective function is then formulated based on these criteria. Kumaradas and Sherar [15] proposed to simultaneously minimize the fraction of damaged volume in the normal tissue and to maximize that of the tumor. This double objective function problem was converted to a single scalar variable using those two factors representing the relative importance of the two minimization criteria. Siauve et al. [16] maximized the ratio of the total heat absorbed by the tumor to that by the healthy tissue using a genetic algorithm method. In a more recent study, Bagaria and Johnson [17] performed optimization calculations to match the obtained temperature profile with an ideal temperature distribution in both the tumor and normal tissue. The objective function consisted of two weighted penalty functions. The first is related to the square error between the tumor temperatures and the necrosis temperature; and the second is related to the square error between the normal tissue temperature and the normal body temperature. It was also found that the optimization of the SAR distribution significantly enhanced the thermal damage to the tumor. However, the optimization was performed based on a quadratic decaying function proposed for the SAR in the tumor region during a magnetic nanoparticle hyperthermia. This expression was not based on any physical measured entities and no explanation was provided for the choice of the function. To the authors’ best knowledge, a general optimization of three-dimensional heating pattern and multi-site injection strategy for clinical applications has not been conducted for nanoparticle hyperthermia.

The objective of the current study is to develop a general algorithm for physicians to design an optimized multi-site injection strategy. The Pennes bioheat equation is first used to determine the thermally affection region ( $d_{43}$ ) based on a single injection in an infinitely large tissue domain. Given an arbitrary irregularly shaped tumor embedded in a normal tissue region, an initial guess of the locations of multiple injection sites and the associated SAR distribution can be made based on the results from the Pennes equation. The Pennes bioheat equation is then applied to determine the temperature fields in both the tumor and normal tissue for the implemented injection sites and heat source. Through a systematic optimization strategy, we obtain the SAR distribution-related parameters (the values for  $A$  and  $r_0$ ) for each injection site to maximize the total

volume of tumor above 43°C while minimizing the temperature elevations in the normal tissue. Since the SAR distribution should be correlated with the injection amount and injection rate at an individual injection site, the appropriate injection strategy can be thus obtained in future clinical studies.

**Model formulation**

The model proposed in this study will give guidance to physicians to select appropriate ferrofluid amounts, injection rates and initial ferrofluid concentration in a multi-site injection procedure of magnetic nanoparticle hyperthermia. As mentioned in the introduction, these entities are directly related to the parameters in the SAR expression  $A_i$  and  $r_{0,i}$  where  $i$  is the index of an injection site. In fact, we found in our previous experimental study in agarose gel [13] that for a given gel concentration, a high injection rate leads to a higher concentration of the nanoparticles close to the injection site. Therefore, a higher injection rate would more likely lead to a bigger value of  $A$  and a smaller  $r_0$ . In our in vivo experiment [14], it was shown that a bigger amount of injected nanoparticles leads to a higher concentration next to the injection site, hence, to a bigger  $A$  and a smaller  $r_0$ . The Pennes bioheat equation [18] is applied first to provide the initial guess of the SAR expression at each injection site for the optimization. We then develop a tumor model with an irregular shape and embedded in a tissue region. Using the Pennes bioheat equation to determine the 3-D temperature fields in both the tumor and normal tissue, we optimize a proposed objective function so that the damage to the tumor is maximized while it is minimized in the normal surrounding tissue. Currently, the Levenberg-Marquardt method, the conjugate directions methods and Simplex methods are techniques for solving optimization problems [19, 20]. In complicated heat transfer optimization problems, the algorithm requires a numerical method for each iteration to determine the temperature field. In such cases, the gradients of the objective function are unknown. Hence, the derivative-free Nelder Mead Simplex method is implemented in this study as an optimization method since its algorithm does not involve computing the gradients of any function.

*The initial guess of the  $d_{T_n}$  using the Pennes bioheat equation*

One of the limitations of the Nelder-Mead Simplex method is its sensitivity to the choice of the initial parameters. To provide a reasonable initial guess to the SAR distribution, we first employ the Pennes bioheat equation to simulate the thermally affected

region induced by a single injection of nanoparticles. The thermally affected region is defined as a region which temperatures are above certain threshold,  $T_n$ . The SAR expression is assumed following the spherical Gaussian’s distribution,

$$SAR = A \cdot e^{-r^2/r_0^2} \tag{1}$$

where  $A$  is the maximum value of the SAR at the center of the injection and  $r_0$  is a dimensional radius that is associated with how far from the injection site the heating is affecting the tissue. Substituting the SAR expression into the steady state Pennes bioheat equation, one obtains,

$$\frac{k}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T}{\partial r} \right) + \rho_b \cdot c_{p,b} \cdot \omega_{b,t} \cdot (T_a - T) + A \cdot e^{-r^2/r_0^2} = 0 \tag{2}$$

where  $k$  is tissue thermal conductivity,  $T$  is temperature,  $\rho_b$ ,  $c_{p,b}$  and  $\omega_{b,t}$  are blood density, specific heat and specific local perfusion rate, respectively, and  $T_a$  is local arterial blood temperature. Note that the temperature field is one-dimensional and only depends on the radial distance from the injection site. It can be solved analytically. The detailed derivation can be found in the Appendix. The following equation gives the steady state temperature distribution in tissue,

$$T(r) = T_a + \frac{A \cdot r_0^3}{2 \cdot r \cdot k \cdot \sqrt{\pi}} \times \int_0^\infty \frac{\beta \cdot e^{-\beta^2/4}}{(\rho_b \cdot c_{p,b} \cdot \omega_{b,t} \cdot (r_0^2/k) + \beta^2)} \cdot \sin\left(\beta \cdot \frac{r}{r_0}\right) \cdot d\beta \tag{3}$$

The temperature field depends on the radial distance, local blood perfusion rate, and the SAR-related parameters. Since the temperature distribution is spherical, one can define an entity  $d_{T_n}$  as the diameter of the thermally affected region, inside which the temperature is higher than or equal to  $T_n$  (Figure 1). For each set of  $A$  and  $r_0$ , and a given blood perfusion rate, a value of  $d_{T_n}$  can be determined based on the calculated temperature distribution. As described later,  $d_{T_n}$  will be used to identify the initial guess for solving the optimization problem. For each value of  $d_{T_n}$  there exist more than one set of  $A$  and  $r_0$ . Hence, to select the initial set, either  $A$  or  $r_0$  should be fixed and the other be calculated, as described later in the section of optimization scheme.

*Mathematical model of the tumor and surrounding tissue*

An irregularly shaped tumor is assumed to be embedded in a large cubical domain of normal tissue as shown in Figure 2 and Figure 3. The size of

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the cubical domain was chosen large enough so that the boundary effects on the temperature distribution are negligible. The thermal properties of the tumor, tissue and blood are assumed to be constant. The three-dimensional Pennes bioheat equation [18] is given in Equation 4 for both the normal tissue and tumor temperature, respectively.

$$k \cdot \left( \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) + \rho_b \cdot c_{p,b} \cdot \omega_b \cdot (T_a - T) + q_m + \sum_{i=1}^M A_i \cdot e^{-r_i^2/r_{0,i}^2} = 0 \tag{4}$$

where  $T$  is either tissue or tumor temperature,  $k$  is the thermal conductivity of the tissue and tumor,  $\rho_b$  and  $c_{p,b}$  are the blood density and the blood

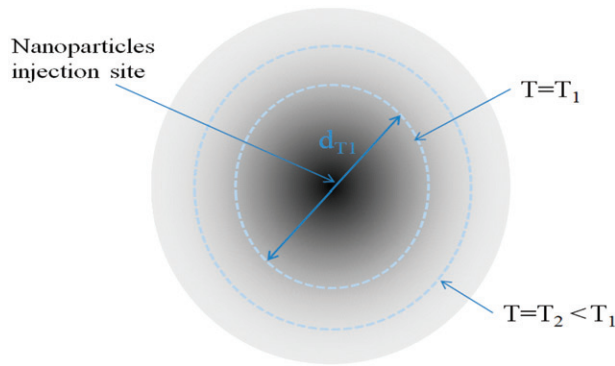


Figure 1. A schematic diagram of the temperature distribution with a single injection in an infinitely large region.

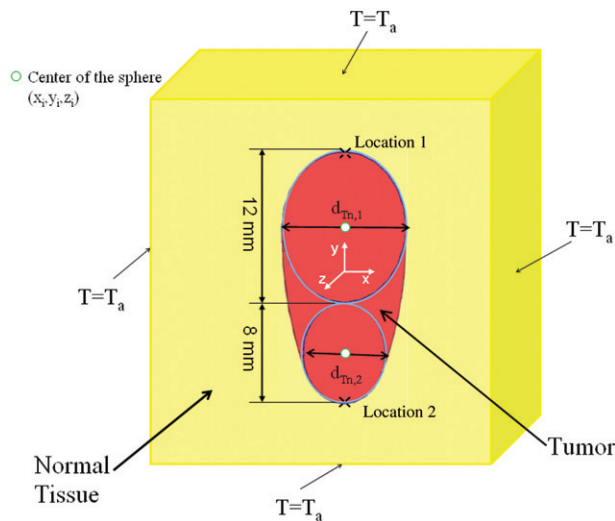


Figure 2. A proposed tumor model with a simple geometry. The tumor is embedded in a cubic healthy tissue region. Two spheres are drawn to match the tumor

specific heat in both tumor and normal tissue,  $\omega_b$  is the local blood perfusion rates in the normal tissue and tumor. The last term in Equation 4 represents the heat sources generated in the vicinity of the nanoparticle injection sites and the nanoparticle heating occurs mainly in the tumor region.  $M$  is the number of injections and  $r_i$  is the radial distance from injection site  $i$  and is given by Equation 5,

$$r_i^2 = (x - x_i)^2 + (y - y_i)^2 + (z - z_i)^2 \quad i = 1, 2, \dots, M \tag{5}$$

where  $x_i$ ,  $y_i$  and  $z_i$  are the Cartesian coordinates of injection site  $i$ . The boundary conditions for this problem are applied so that the temperatures at the outer border of the normal tissue are equal to  $T_a$ . The heat generation term in Equation 4 was programmed and implemented in the optimization algorithm using the Matlab interface of Comsol® Multiphysics.

**Optimization scheme**

*Objective function*

In any hyperthermia treatment for tumor, one would ensure a complete necrosis of the tumor region while protecting the normal tissue. Such task is fulfilled when the temperature inside the tumor is higher than or equal to  $T_n$  while the temperature in the normal tissue region is less than  $T_n$ . Furthermore, to ensure that the whole tumor acquired the temperature  $T_n$ , the boundary separating the tumor from its surrounding tissue also has to achieve the same temperature. In other words, the volume of the region inside the tumor and the tumor boundary area having a temperature higher than  $T_n$

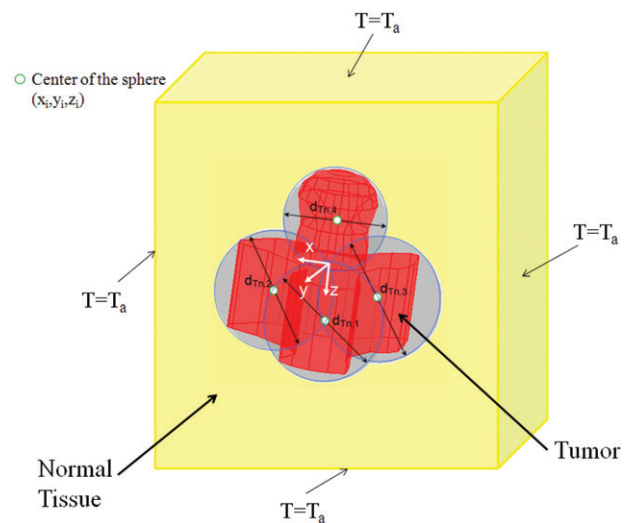


Figure 3. A tumor with an irregular geometry is embedded in a cubic tissue region. Four spheres located at different locations are needed to match the tumor shape.

should be maximized. We first define two percentages  $R_1$  and  $R_2$  as:

$$R_1 = \frac{V_{\text{tumor}, T \geq T_n}}{V_{\text{tumor}, \text{total}}} \cdot 100 \quad (6a)$$

$$R_2 = \frac{S_{\text{tumor boundary}, T \geq T_n}}{S_{\text{tumor boundary}, \text{total}}} \cdot 100 \quad (6b)$$

$R_1$  is the volume percentage of the tumor region with  $T \geq T_n$  over the entire tumor region and  $R_2$  is the area percentage of the tumor boundary where  $T \geq T_n$  over the entire tumor surface. During optimization, the magnitude of the nanoparticle heating in the tumor can be increased so that both  $R_1$  and  $R_2$  can reach 100%. However, the constraint here is minimizing the damage to the surrounding normal tissue. One way to address the constraint is to introduce a third parameter  $E$ , which represents the deviation of the tumor boundary temperature from  $T_n$  [17]. The dimensionless parameter  $E$  is defined as the summation of the square error between the temperature of the tumor boundary and  $T_n$ .

$$E = \frac{1}{(T_a - T_n)^2} \cdot \sum_{j=1}^N (T_j - T_n)^2 \quad (7)$$

where  $T_j$  is the temperature at the tumor boundary (node  $j$ ). Note that all three parameters are dimensionless. Based on the above considerations, the objective is to maximize both  $R_1$  and  $R_2$ , as well as to minimize  $E$ . The expected temperature field after optimization would yield a minimal deviation of the isotherm line corresponding to  $T_n$  from the tumor boundary, as shown in Figure 4. We notice

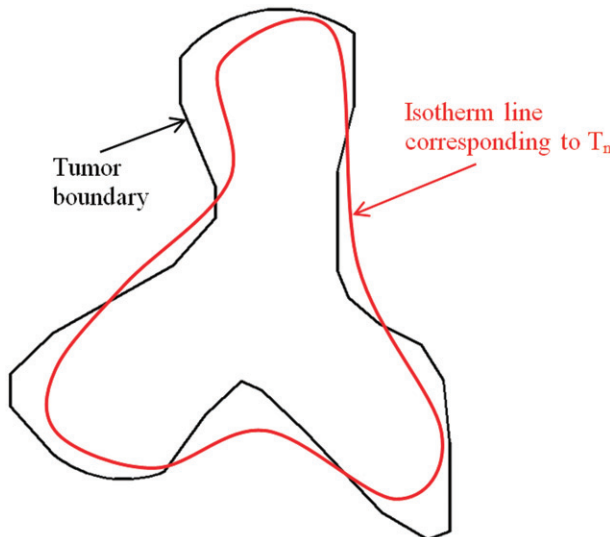


Figure 4. Schematic diagram of the tumor boundary and the isotherm line corresponding to  $T_n$ .

that our optimization problem involves both maximizing and/or minimizing different quantities simultaneously [15, 16]. This task should be expressed in the objective function proposed in Equation 8.

$$f = \frac{E}{R_1 \cdot R_2} \quad (8)$$

We use the Nelder-Mead Simplex method [21] that can handle non-linearities and discontinuities to minimize the objective function  $f$  and to optimize both  $A_i$  and  $r_{0,i}$ .

The effectiveness of the nanoparticle hyperthermia is evaluated by two parameters, one is the  $T_{90, \text{tumor}}$  in the tumor region and the other is the  $T_{10, \text{tissue}}$  value in the surrounding tissue region.  $T_p$  is defined as the temperature exceeded by  $p\%$  of a certain region [22]. Both inequalities  $T_{90, \text{tumor}} \geq T_n$  and  $T_{10, \text{tissue}} \leq T_n$  suggest that the treatment will lead to more than 90% of the tumor region reaching the threshold temperature of  $T_n$ , while less than 10% of the normal tissue region is subjected to thermal damage.

#### The optimization procedure

The optimization procedures run as follows.

1. Based on the geometry of the irregularly shaped tumor, several spheres are drawn in the tumor. In the current study, we assume that the centers of the spheres are pre-selected and they are the locations of the injection sites such that the spheres cover almost all the tumor periphery as shown in Figure 2 and Figure 3. The diameter of each sphere,  $d_{T_n, \text{init}}$  is measured.
2. Based on the solution of the Pennes bioheat equation on an infinite domain (see the Appendix), suitable sets of  $A_{\text{init}}$  and  $r_{0, \text{init}}$  are calculated based on the measured  $d_{T_n, \text{init}}$ . Since there are more than one set of  $A$  and  $r_0$  available for any  $d_{T_n}$ , we select  $r_{0, \text{init}}$  as equal to  $r_{0, \text{init}} = (d_{T_n, \text{init}}/2)/\sqrt{3}$ . This is based on the consideration that 95% of the nanoparticles are injected within this spherical domain. Once  $r_0$  is known, the initial value of  $A$  can then be determined based on Figure 5.
3. Those  $A$  and  $r_0$  values for different injection sites are used as the initial guess for the optimization problem and the Nelder-Mead Simplex method is used to minimize  $f$ .

#### Tumor geometry and numerical method

We performed case studies using two geometric shapes of tumor as shown in Figures 2 and 3 to demonstrate the feasibility of the optimization algorithm. The first geometry is axis-symmetric. The second shape was randomly modeled on Pro

Engineer Wildfire 2.0 and then imported to the numerical method software Comsol® Multiphysics. Both tumors have an overall longitudinal size of 20 mm. The simulation domain of the tissue which is a 4096 cm<sup>3</sup> cube is very large to ensure that the boundary condition on the surface is equal to 37°C. However, we calculated the  $T_{10,tissue}$  based on a much smaller cubical domain surrounding the tumor and having a length only twice the tumor's average dimensions. The volumes of the tissue region for the  $T_{10,tissue}$  are different for both cases since they are based on the tumor average dimensions. The volume of the first and second tumors is 1.55 and 2.14 cm<sup>3</sup> respectively. The differential equations were solved using the finite element solver of Comsol® Multiphysics 3.4. The mesh was finer inside the tumor with a growth factor of 1.1 starting from the tumor boundary toward the normal tissue. The total number of the quadratic elements and degrees of freedom were around 166,000 and 226,000, respectively, for both tumor models. The mesh sensitivity was checked by multiplying the total number of the elements by four and finding that the global result differed by less than 1%. The optimization computations were performed at both variable and function tolerance of  $10^{-4}$ .

## Results

All the computations in this study are performed for  $T_n = 43^\circ\text{C}$ , which is considered as a threshold temperature to induce cell necrosis within a reasonable time duration. The values of the thermal conductivities for the normal tissue  $k_{tissue}$  and for the tumor  $k_{tumor}$  are chosen to be 0.5 and 0.55 W/m°C, respectively, and the blood density  $\rho_b$  and specific

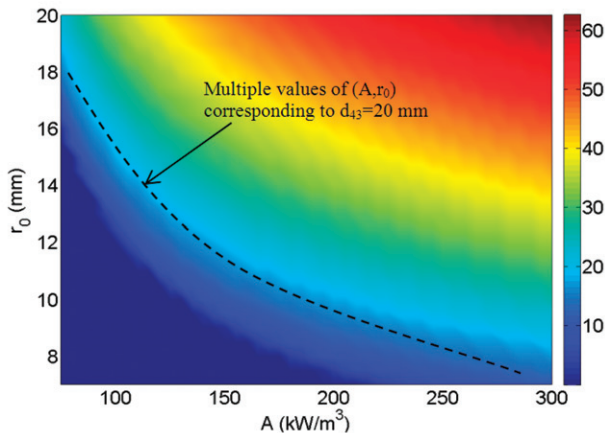


Figure 5. A contour plot of the diameter  $d_{Tn}$  as a function of  $A$  and  $r_0$ . The calculation is based on the simulation in the infinite domain with single injection.  $T_n = 43^\circ\text{C}$  and  $\omega_{b,t} = 1.25 \cdot 10^{-4} \text{ s}^{-1}$ .

heat  $c_{p,b}$  are assumed to be 1000 Kg/m<sup>3</sup> and 4200 J/Kg°C respectively [17]. The initial guess is calculated based on a specific blood perfusion rate in the tumor  $\omega_{b,t}$ . For  $\omega_{b,t} = 1.25 \cdot 10^{-3} \text{ s}^{-1}$  [17], we calculate  $d_{Tn}$  for multiple values of  $A$  and  $r_0$ . The results are shown in Figure 5. Clearly, higher values of  $A$  and  $r_0$  result in a bigger temperature region higher than  $43^\circ\text{C}$  and lead to a bigger value of  $d_{43}$ . For the same  $d_{43}$ , more than one sets of  $A$  and  $r_0$  are available (see Figure 5).

The optimization algorithm was applied to the first (simple) tumor. Since the local blood perfusion rate of tumor may be quite different from that of normal tissue [23], we selected the blood perfusion rate in tissue as  $\omega = 5 \cdot 10^{-4} \text{ s}^{-1}$  [14, 24]. Based on the tumor geometry, two injection sites are selected and the initial guess of  $A$  and  $r_0$  were estimated (Table I). The optimization algorithm converged after approximately 60 iterations. For the first tumor, one notices that the initial guess leads to a  $T_{10,tissue}$  that is greater than  $43^\circ\text{C}$ , even if the initial  $T_{90,tumor}$  has satisfied the design objective. The optimization has resulted in a much smaller value of  $T_{10,tissue}$  ( $41.79^\circ\text{C}$ ), which is about two degrees less than its initial value. The optimization result is further illustrated by the color contours in Figure 6. The boundary of the tumor is represented by the solid line while the isotherm of  $43^\circ\text{C}$  can be seen by the light blue line. Small deviation between the tumor surface and the isotherm is evident in Figure 6, especially in the middle portion of the tumor where the tumor surface deviates from the spherical shape.

For the first geometry, analyses were performed to check for the sensitivity of the optimization result to the initial guess of  $A$  and  $r_0$ . The initial values of  $r_{0,1}$  and  $A_1$  were either increased or decreased by 5%. Figure 7 shows that  $r_0$  and  $A$  converge to the same values when the initial guesses vary  $\pm 5\%$ . Bigger variations in the initial guesses did not result in the same optimum solution. As shown in Table II, changing the initial values by 10% or 20% will result in changes in  $A$  and  $r_0$ , however, the optimization objective of  $T_{90,tumor} > 43^\circ\text{C}$  and  $T_{10,tissue} < 43^\circ\text{C}$  is still satisfied.

As stated in the model formulation, the injection sites are pre-selected based on the geometry of

Table I. Summary of the initial and optimized values for the first tumor geometry shown in Figure 2 when  $\omega = 5 \cdot 10^{-4} \text{ s}^{-1}$ ,  $\omega_t = 1.25 \cdot 10^{-3} \text{ s}^{-1}$ .

		$d_{43}$	$r_0$	$A$	$T_{90,tumor}$	$T_{10,tissue}$
	Site	(mm)	(mm)	(MW/m <sup>3</sup> )	(°C)	(°C)
Initial values	1	12	3.5	1.79	45.51	43.44
	2	8	2.3	3.31		
Final values	1	11.4	3.1	1.91	43.43	41.79
	2	7.6	2.0	3.59		

the tumor. A poorly selected injection site may not yield a design which satisfies both  $T_{90,\text{tumor}} > T_n$  and  $T_{10,\text{tissue}} < T_n$ . The sensitivity of the optimization algorithm to the pre-determined injection sites was

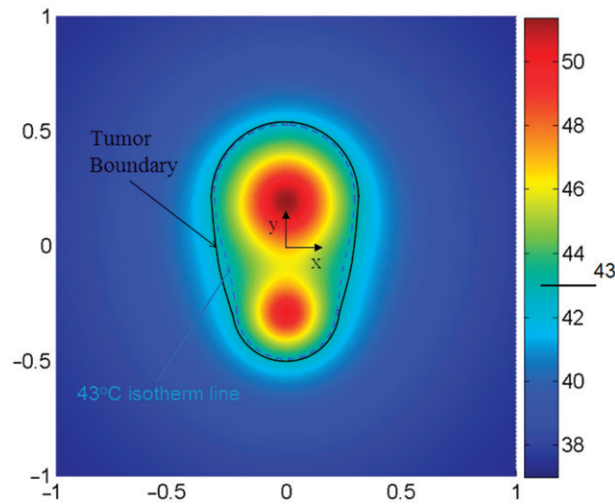


Figure 6. Temperature contours in the simple tumor geometries after optimization. The tumor boundary is represented by the black solid line and the isotherm of 43°C is by the light blue dash line. Note that the  $z$ -direction is perpendicular to the 2-D  $x$ - $y$  plane. The coordinates are normalized by 20 mm which is the longitudinal length of the tumor.

evaluated by changing the location of the sites by a small variation along the axis of the tumor. We consider two situations: either the injection sites are moved away from each other or towards each other. The results are expressed in terms of the  $T_{90,\text{tumor}}$  and  $T_{10,\text{tissue}}$  in Table III. Our results have demonstrated that moving the injection locations away or close to each other by 2.5% of the total length of the tumor leads to a minor increase in  $T_{90,\text{tumor}}$  with no significant change in  $T_{10,\text{tissue}}$ . However, moving the injection sites by more than 10% fails to guarantee a  $T_{90,\text{tumor}}$  bigger than 43°C, as shown by the bold columns in Table III.

We also test the sensitivity of the optimum values of  $A$  and  $r_0$  to changes in the blood perfusion rate in the tumor. The resulting optimum values of  $A$  can be affected significantly as shown in Table IV. Our results have demonstrated that an increase of 10% or 50% in the tumor blood perfusion rate results in an increase in the value of  $A$  by 3% and 16% respectively while satisfying almost the same treatment goals. We note that the optimization algorithm still works fine with different blood perfusion rates. This is expected since local blood perfusion rate has been demonstrated to play a significant role in determining the local temperature field in tumor. The large blood perfusion rate acts as a heat sink and carries away the heat induced by the nanoparticle heat generation.

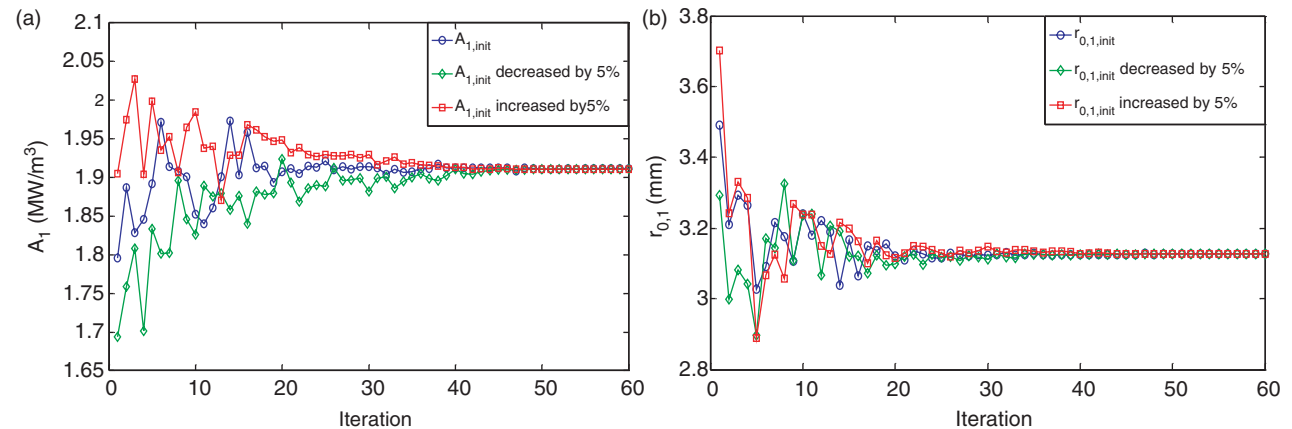


Figure 7. Convergence curves for  $A_1$  (a) and  $r_{0,1}$  (b) and the sensitivity of the convergence to a 5% change in the initial values.

Table II. The effect of changing the initial values of  $A$  and  $r_0$  by 10% and 20% on the optimum values and on the  $T_{90,\text{tumor}}$  and  $T_{10,\text{tissue}}$ .

	Initial values unchanged	$A_{1,\text{init}}$ and $r_{0,1,\text{init}}$ changed by 10%	$A_{1,\text{init}}$ and $r_{0,1,\text{init}}$ changed by 20%
$A_1$ (MW/m <sup>3</sup> )	1.91	1.76	1.21
$r_{0,1}$ (mm)	3.1	3.2	3.6
$T_{90,\text{tumor}}$ (°C)	43.43	43.42	43.41
$T_{10,\text{tissue}}$ (°C)	41.79	41.79	41.84



Table III. The effect of varying the injection sites by a small distance on the optimization results for the first tumor geometry.

Non-dimensional distance	Injection sites moved away from each other by:		Originally selected injection locations	Injection sites moved towards each other by:	
	10%	2.5%		2.5%	10%
$T_{90,tumor}$ ( $^{\circ}\text{C}$ )	42.88	<b>43.27</b>	43.43	43.37	<b>42.63</b>
$T_{10,tissue}$ ( $^{\circ}\text{C}$ )	42.39	<b>41.81</b>	41.79	41.76	<b>41.5</b>

Table IV. The effect of increasing the tumor blood perfusion rate on the optimization results for the first tumor geometry.

	Site	$\omega_t =$	$\omega_t$ increased	$\omega_t$ increased
		$1.25 \cdot 10^{-3} \text{ s}^{-1}$	by 10%	by 50%
$A$ ( $\text{MW/m}^3$ )	1	1.91	1.96	2.22
	2	3.59	3.65	3.98
$r_0$ (mm)	1	3.1	3.1	3.0
	2	2.0	2.0	1.9
$T_{90,tumor}$ ( $^{\circ}\text{C}$ )		43.43	43.43	43.44
$T_{10,tissue}$ ( $^{\circ}\text{C}$ )		41.79	41.79	41.8

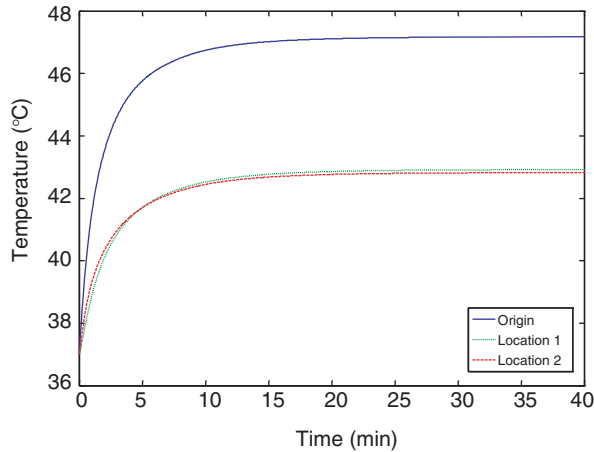


Figure 8. The transient temperature elevation at three locations in the first tumor geometry using the optimized heating pattern.

In this paper, we perform the optimization process in steady state to determine the heating pattern by nanoparticles due to the restriction of the available computational resources. Although simulating the transient process of the temperature distribution in tumor is not the focus of the current paper, we performed simulations to predict the transient temperature elevations in the tumor using the optimized heating pattern for the first geometry. As shown in Figure 8, temperature distribution inside the tumor is plotted for various time instants. It takes less than 30 minutes to establish a steady state. For clinical

applications, the appropriate treatment (heating) time should be determined based on the equivalent  $EM_{43}$  for the entire tumor region. Usually an  $EM_{43}$  is longer than 90 minutes to achieve complete tumor necrosis.

To test whether the developed algorithm is equally capable of handling 3-D tumor with complicated geometry, we tested the algorithm on the second tumor. Table V gives the selected coordinators of the four injection sites, as well as the initial and optimized values of  $r_0$  and  $A$ . The convergence of the program took about 150 iterations (Figure 9) and the computation time was approximately 24 h due to the complex tumor geometry and four injection sites. The effectiveness of the algorithm is illustrated by Table VI, where the initial and final  $T_{90,tumor}$  and  $T_{10,tissue}$  are given. The success of the algorithm is evident by the converged values of  $A$  and  $r_0$  at each injection site as well as the resultant  $T_{90,tumor}$  ( $43.21^{\circ}\text{C}$ ) and  $T_{10,tissue}$  ( $41.76^{\circ}\text{C}$ ).

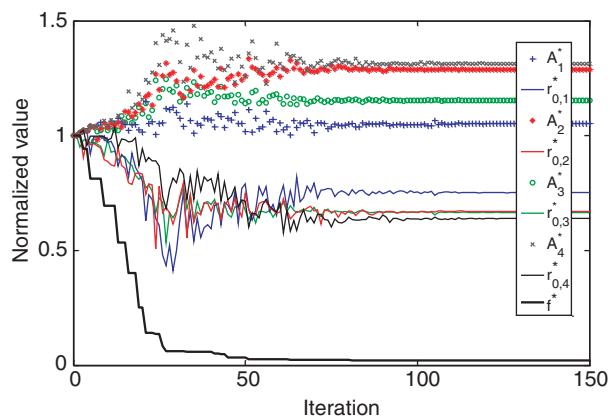
Figure 10 gives the initial and optimized temperature fields of the second tumor geometry. Overheating in the healthy tissue is evident in the initial temperature contours. After the optimization, only a small region ( $<3\%$ ) of the healthy tissue is found above  $43^{\circ}\text{C}$  while more than 90.2% of the tumor region has been elevated to bigger than  $43^{\circ}\text{C}$ . Deviation between the  $43^{\circ}\text{C}$  isotherm and the tumor boundary can also be seen in Figure 10(b).

## Discussion and conclusion

In this paper we develop an algorithm to optimize the heat absorption distribution for multi-site injections in magnetic nanoparticle hyperthermia when a single injection is not sufficient to cover the entire irregular-shaped tumor. The output of this study is an optimized SAR distribution in terms of  $A$  and  $r_0$ . The effectiveness of the optimization algorithm is evaluated by  $T_{90,tumor}$  and  $T_{10,tissue}$ . The concept of  $T_{90}$ ,  $T_{50}$  and  $T_{10}$  has been widely used and reviewed in hyperthermia treatment in both tumor and tissue [25–30]. Theoretically, one would expect that  $T_{90} < T_{50} < T_{10}$  in any tumor or normal tissue

Table V. Summary of the injection site locations and the initial and optimized values of  $d_{Tn}$ ,  $A$  and  $r_0$  for the second tumor geometry.

	1	2	3	4
$x$ (mm)	1.56	-6.48	4.3	0
$y$ (mm)	9.02	-2.6	-6.42	0
$z$ (mm)	4.5	4.5	4.5	-6.7
$A_{init}$ (MW/m <sup>3</sup> )	1.63	1.7	1.6	2.1
$r_{0,init}$ (mm)	3.7	3.59	3.78	3.11
$d_{43,init}$ (mm)	12.83	12.46	13.11	10.8
$A$ (MW/m <sup>3</sup> )	1.87	2.076	1.75	2.76
$r_0$ (mm)	2.71	2.41	2.67	2.07
$d_{43}$ (mm)	9.4	8.8	9.2	7.2


 Figure 9. The convergence curves of  $A$  (symbols),  $r_0$  (lines) and the objection function  $f$  (heavy black line) obtained by the Nelder-Mead Simplex method for the second tumor geometry.  $A$  and  $r_0$  were normalized by their initial guess.

region during hyperthermia. For instance, if these three temperatures are comparable to each other, the temperature distribution is almost uniform. However, this is not the case in many hyperthermia treatment situations where it is almost impossible to obtain a uniform high temperature in the tumor and a sudden temperature decrease to the normal tissue. Therefore, we chose to implement  $T_{90}$  and  $T_{10}$  to the tumor and tissue, respectively. Since we have to achieve thermal damage to the tumor, we should ensure a  $T_{90}$  bigger than 43°C. The opposite should take place in the healthy tissue meaning that  $T_{10}$  should be less than 43°C. In this paper, we present two case studies to apply the algorithm we developed.

The focus of this paper is to demonstrate the feasibility of the optimization algorithm. We are confident that the algorithm will work if another treatment parameter such as  $T_{5,tissue}$  or  $T_{1,tissue}$  is chosen. The choice of  $T_{10}$ ,  $T_5$  or  $T_1$  in the healthy tissue region is totally dependent on the choice of clinicians, since it depends on the tumor location,

 Table VI. Summary of the optimized values of the percentages  $R_1$  and  $R_2$ , the normalized square error and the thermal damage indicators  $T_{90}$  and  $T_{10}$  in the tumor and normal tissue for the second tumor geometry.

	Initial	Final
$R_1$ (%)	100	90.7
$R_2$ (%)	100	69.36
$E/N$	2.53	0.029
$f$	2.53	0.047
$T_{90,tumor}$ (°C)	52.22	43.21
$T_{10,tissue}$ (°C)	48.93	41.76

shape and size. The clinicians need to specify a margin of the healthy tissue that has to be protected during the treatment. In future theoretical and clinical studies, one could investigate the effect of the size of the healthy tissue domain surrounding the tumor on the outcome of the optimization algorithm, i.e the value of  $T_{10,tissue}$ . If the defined tissue region is very small to result in an optimized  $T_{10,tissue}$  higher than 43°C, it only suggests that more injection sites may be needed and/or the injection locations are re-adjusted. In this paper, we present two case studies to demonstrate that the algorithm we developed works in typical clinical situations. The theoretical model has the potential of simulating and optimizing realistic situations of magnetic nanoparticle hyperthermia. Unlike previous theoretical studies [17] where the SAR expression was not based on experimental measurements, our model is more realistic because the SAR expressions were obtained from experimental results in tissue-equivalent agarose gel and animal tissue. Furthermore, our approach is more general since it does not rely on weighting and relative importance factors in the choice of the objective function. To our knowledge, the current study is the first attempt to simulate and generally optimize the three-dimensional heating pattern to determine optimal injection parameters in clinical applications of magnetic nanoparticle hyperthermia.

On the other hand, the proposed optimization model has limitations. Uncertainties concerning the tumor location and injection sites are inevitable in clinical applications. The accuracy of the injection sites and sizes of the spheres depends on the accuracy of the modern imaging techniques such as MRI or CT. During clinical study, the optimization algorithm can be applied to a tumor with similar shape, however, it may be slightly bigger or smaller within the measurement margin of the imaging technique. We can provide the clinicians with multiple treatment protocols reflecting the inaccuracy of the size and/or shape of the tumor. However, if the resulting temperature outcomes are significant, the clinicians

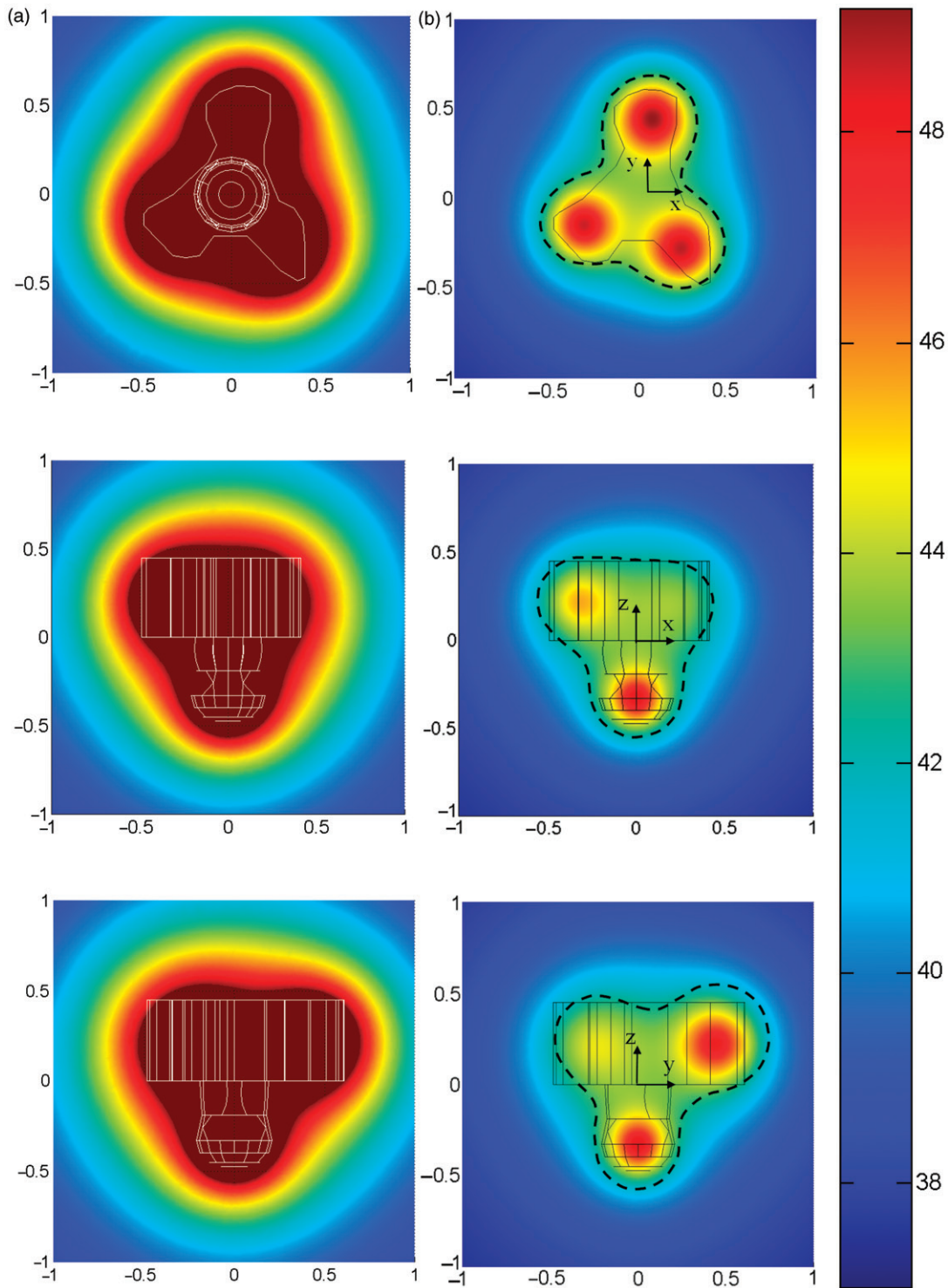


Figure 10. The optimum temperature distributions in the second tumor geometry on three cross-sectional planes passing through the injection sites: (a) before optimization; (b) after optimization (the dashed black line represents the 43°C isotherm). Note that the third coordinate is perpendicular to the 2-D plane. The coordinates are normalized by 20 mm which is the longitudinal length of the tumor.

may utilize another imaging technique which has a better measurement accuracy.

In this study, the locations of the injection sites are pre-selected. The injection sites and the initial sphere

diameters  $d_{43}$  are determined in a way so that the thermally affected region maximally matches the actual geometry of the tumor. The diameter of the sphere is selected so that 95% of the injected

nanoparticles are located within the sphere. If the optimization result does not satisfy the treatment goal ( $T_{90, \text{tumor}} > 43^\circ\text{C}$  and  $T_{10, \text{tissue}} < 43^\circ\text{C}$ ), it suggests that more injection sites are needed than the original sites. As shown in the results, the pre-selection of the locations may or may not lead to an ideal temperature distribution or to a design which satisfies overheating the tumor while protecting the normal tissue, even if the computational algorithm converges to the optimal result. If the result fails to satisfy the design goals, one needs to examine the obtained temperature contours in both the tumor and tissue regions and adjust the injection locations and/or the number of the injection sites to better match the tumor geometry. Moreover, more accurate results can be obtained by considering the three coordinates of each injection site as optimization variables. This approach makes the algorithm computationally expensive since the three spatial coordinates of each injection site are added to the two SAR parameters  $A$  and  $r_0$ . Therefore, the number of variables becomes  $5M$ , where  $M$  is the number of injection sites whereas in the current study it is equal to  $2M$ . The developed algorithm can be modified to also optimize the injection locations in the future when computational resources are available.

One notes that the optimization procedures may yield to more than one set of the SAR parameters. Since both  $A$  and  $r_0$  contribute equally to the heat deposition in the tumor, one has the flexibility to select different combinations of the parameters. However, there is a limit on how high a maximum SAR at the injection site can be achieved. It is well known that the maximum SAR is affected by the particle properties, available magnetic field strength and frequency, as well as the initial ferrofluid concentration, injection amount and injection rate [2, 3, 13, 14]. In fact, in our ongoing research of nanoparticle hyperthermia, we will initially manipulate the injection volume and rate. The effects of other parameters dependent on the equipment and the available ferrofluids in the market will be investigated later. Future experimental studies performed on tumors of various types in animal and clinical studies are crucial to understanding the limit of SAR strength. Nevertheless, our results have shown that a small variation from the initial guess would converge to the same optimization results of the SAR distributions.

We envisage that the current study has great potential in the clinical study of nanoparticle hyperthermia. Before this can be used in clinical applications, extensive experimental studies are required to generate a database that relates the injection parameters, such as quantity of nanoparticles and injection rate, to the SAR distribution parameters  $A$  and  $r_0$ . It is expected that the clinician

will first scan the geometry of the tumor and tissue and measure the local blood perfusion rates of both regions using MRI or other non-invasive imaging techniques. The injection sites are determined based on the tumor geometry. Our optimization algorithm is then implemented to the tumor and its surrounding tissue to determine the optimal SAR parameters of all injection sites. Three-dimensional temperature contours can be plotted and reviewed to see whether there is any thermally sensitive region exposed to excessive heating. After that the clinician will consult the database and find the suitable and realistic injection amount and injection rate. When selecting the injection amount and injection rate, one has to keep in mind how long the patient can tolerate the injection time when the injection rate is slow. Another consideration is the response of local blood perfusion rate to heating during nanoparticle hyperthermia. It is well known that hyperemic response to heating can potentially lower the treatment efficacy due to the relatively cold blood supplied to the tumor region [31]. However, the response depends on the tumor type and size, and the response may be different from tumor to healthy tissue [32]. To provide the optimization algorithm with accurate blood perfusion rates, one has to rely on data of experimental studies on heat response during hyperthermia treatment. The advantage of the algorithm developed in the study is that it can be used to provide different scenarios of optimization plan for different blood perfusion rates. Once the injection strategy of each injection site is determined, the Pennes bioheat equation can be re-run to simulate the three-dimensional temperature field during treatment duration. Note that in most hyperthermia treatment, an equivalent time for  $43^\circ\text{C}$  ( $EM_{43}$ ) is usually used to evaluate whether the heating is sufficient to result in tumor cell necrosis [27]. The simulation of the 3-D transient temperature fields in both tumor and surrounding tissue allows for estimating and mapping the  $EM_{43}$  distribution to further evaluate thermal damage to the targeted region.

In conclusion, our study has demonstrated the feasibility of optimizing heating pattern in tumor treatment using magnetic nanoparticle hyperthermia. The optimization algorithm designed in this study significantly improves the treatment outcome, especially when the tumor is deep seated and irregular in shape. Minor changes in the initial values of the SAR parameters do not affect significantly the optimum solution. It has been shown that the injection site location and number of injections play important roles in determining whether the optimization has achieved the targeted  $T_{90, \text{tumor}}$  and  $T_{10, \text{tissue}}$ . Future experimental studies are needed to relate the SAR parameters to the injection strategy for each injection



site. We believe that the current study has the potential to provide guidance to clinicians in designing effective and safe treatment protocols for cancer patients.

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**Appendix**

We estimate  $d_{T_n}$  by solving the Pennes bioheat equation in spherical coordinates in an infinite domain.

$$\frac{k}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T}{\partial r} \right) + \rho_b \cdot c_{p,b} \cdot \omega_{b,t} \cdot (T_a - T) + A \cdot e^{-r^2/r_0^2} = 0$$

$r \rightarrow 0$ ;  $T$  is finite  
 $r \rightarrow \infty$ ;  $T \rightarrow T_a$  (I)

Dimensionless parameters are introduced as

$$r^* = \frac{r}{r_0}$$

$$T^* = \frac{T - T_a}{T_n - T_a} \quad \text{(II)}$$

The dimensionless equation and boundary conditions become:

$$\frac{k}{r_0^2} \frac{1}{r^{*2}} \frac{d}{dr^*} \left( r^{*2} \frac{dT^*}{dr^*} \right) - \rho_b \cdot c_{p,b} \cdot \omega_{b,t} \cdot T^* + \frac{A \cdot e^{-r^{*2}}}{T_n - T_a} = 0$$

$r^* \rightarrow 0$ ;  $T^*$  is finite  
 $r^* \rightarrow \infty$ ;  $T^* \rightarrow 0$  (III)

We apply the following variable transformation,  $u = T^* \cdot r^*$  [33–35] and the equation changes to:

$$\frac{d^2 u}{dr^{*2}} - \alpha^2 \cdot u + C \cdot r^* \cdot e^{-r^{*2}} = 0$$

$r^* \rightarrow 0$ ;  $u \rightarrow 0$   
 $r^* \rightarrow \infty$ ;  $u \rightarrow 0$  (IVa)

where

$$\alpha^2 = \rho_b \cdot c_{p,b} \cdot \omega_{b,t} \cdot \frac{r_0^2}{k} \quad \text{(IVb)}$$

$$C = \frac{A \cdot r_0^2}{k \cdot (T_n - T_a)} \quad \text{(IVc)}$$

Applying the sine Fourier transforms to the above equation (Equation IVa) and solving for the sine Fourier transform of  $u$  yield:

$$\bar{u} = \frac{C \cdot \sqrt{\pi} \cdot \beta \cdot e^{-\beta^2/4}}{4 \cdot (\alpha^2 + \beta^2)} \quad \text{(V)}$$

where  $\beta$  is the sine Fourier transform variable. Finally, applying the inverse sine transform results in:

$$u = \frac{C}{2 \cdot \sqrt{\pi}} \cdot \int_0^\infty \frac{\beta \cdot e^{-\beta^2/4}}{(\alpha^2 + \beta^2)} \cdot \sin(\beta \cdot r^*) \cdot d\beta \quad \text{(VI)}$$

The 1-D temperature distribution can then be written as

$$T = T_a + \frac{A \cdot r_0^3}{2 \cdot r \cdot k \cdot \sqrt{\pi}} \times \int_0^\infty \frac{\beta \cdot e^{-\beta^2/4}}{(\rho_b \cdot c_{p,b} \cdot \omega_{b,t} \cdot (r_0^2/k) + \beta^2)} \cdot \sin\left(\beta \cdot \frac{r}{r_0}\right) d\beta$$

In this study, the thermally affected region is defined as a spherical region where the temperature is higher than  $T_n$ . The diameter of the sphere  $d_{T_n}$  can be determined by solving the following equation.

$$T_n = T_a + \frac{A \cdot r_0^3}{d_{T_n} \cdot k \cdot \sqrt{\pi}} \times \int_0^\infty \frac{\beta \cdot e^{-\beta^2/4}}{(\rho_b \cdot c_{p,b} \cdot \omega_{b,t} \cdot (r_0^2/k) + \beta^2)} \cdot \sin\left(\beta \cdot \frac{d_{T_n}}{2 \cdot r_0}\right) d\beta \quad \text{(VII)}$$